

THE ROLE OF THERMAL PRESSURE IN JET LAUNCHING

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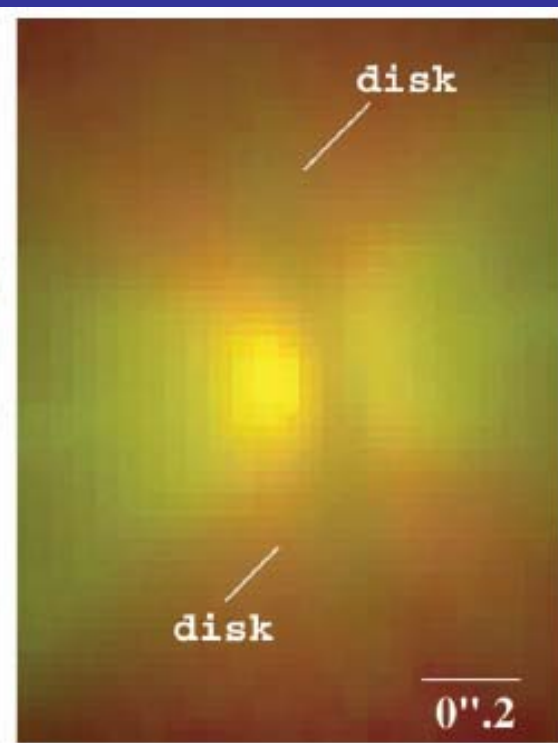
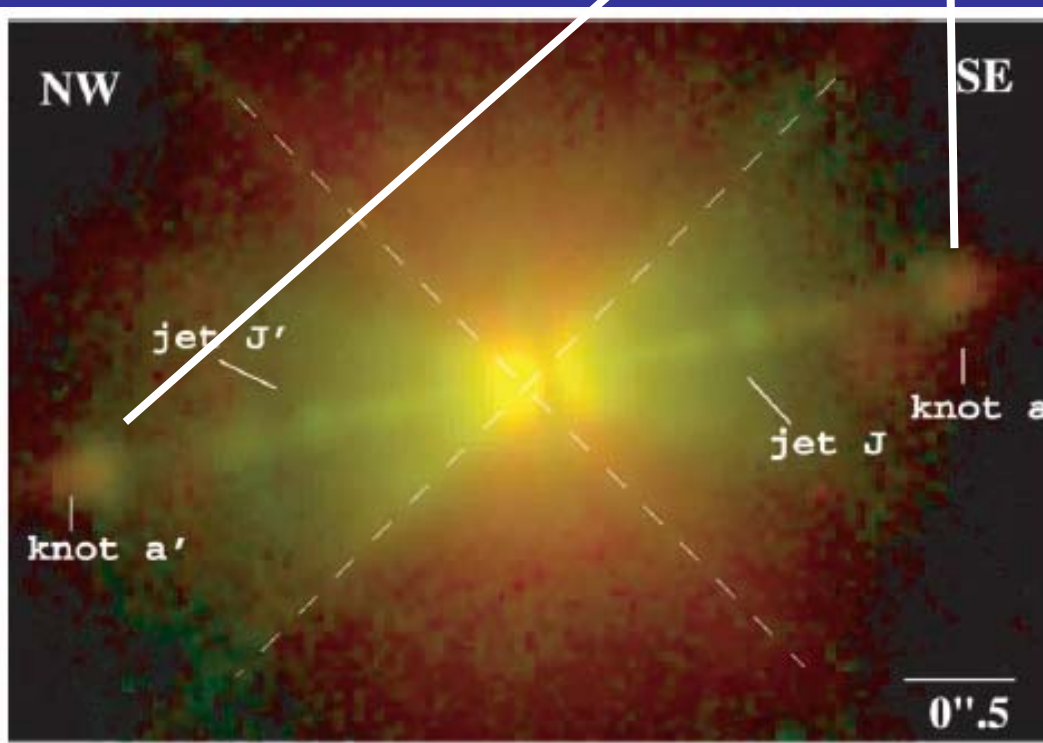
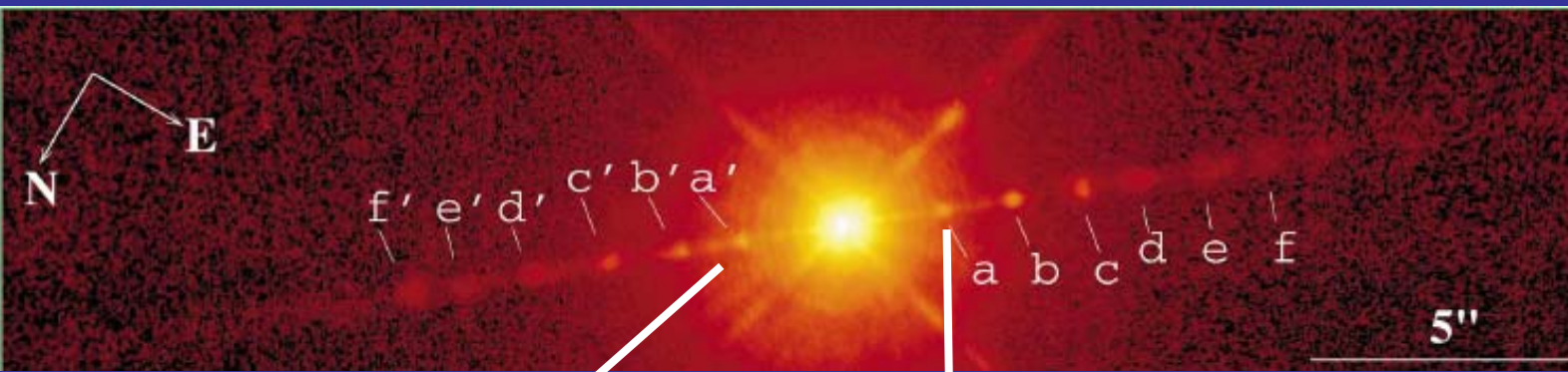
Technion, Israel

MOTIVATION

- ★ In the inner boundary of the disk the gas has the maximum available energy per unit mass.
- ★ Large gradients in one or more of the physical parameters exist. These allow strong dissipation of the disk's kinetic energy.
- ★ Strong dissipation is best done via shocks, which build large pressure gradients, mainly perpendicular to the disk plane.

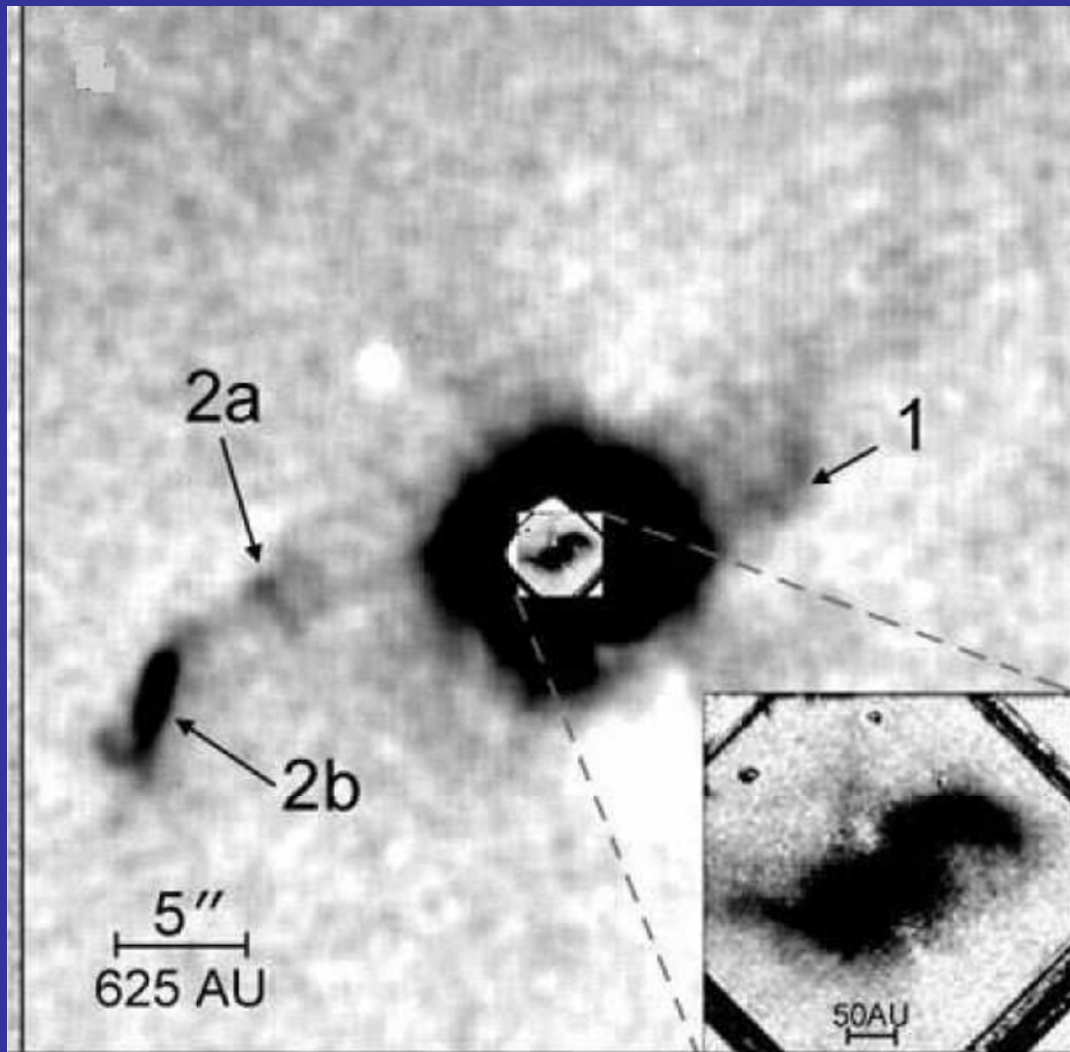
Speed: ~ 200 km/sec

A pair of blobs is ejected every ~ 40 years

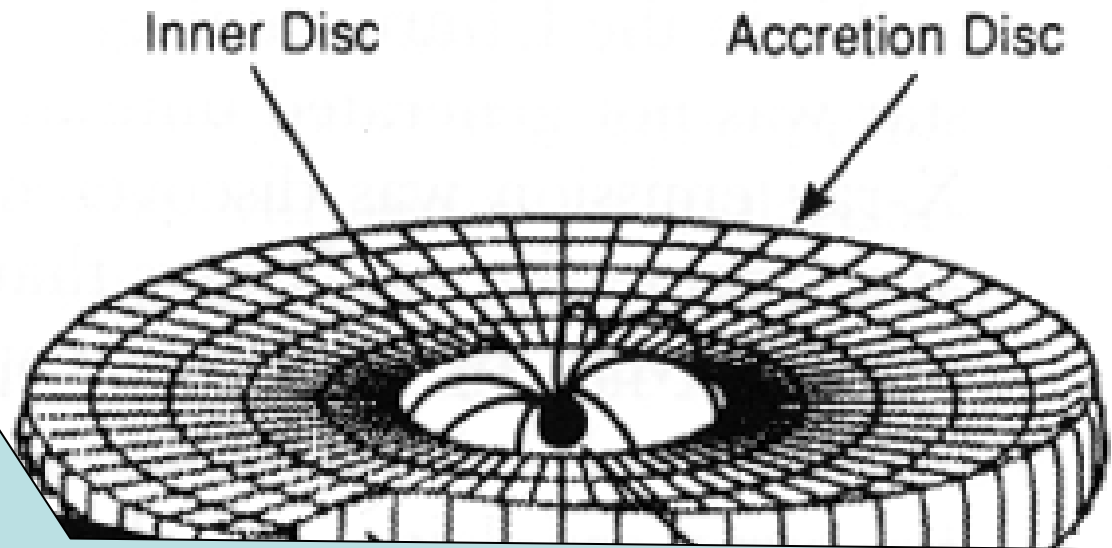


HENIZE 2-90: A planetary nebula (Sahai, R.)

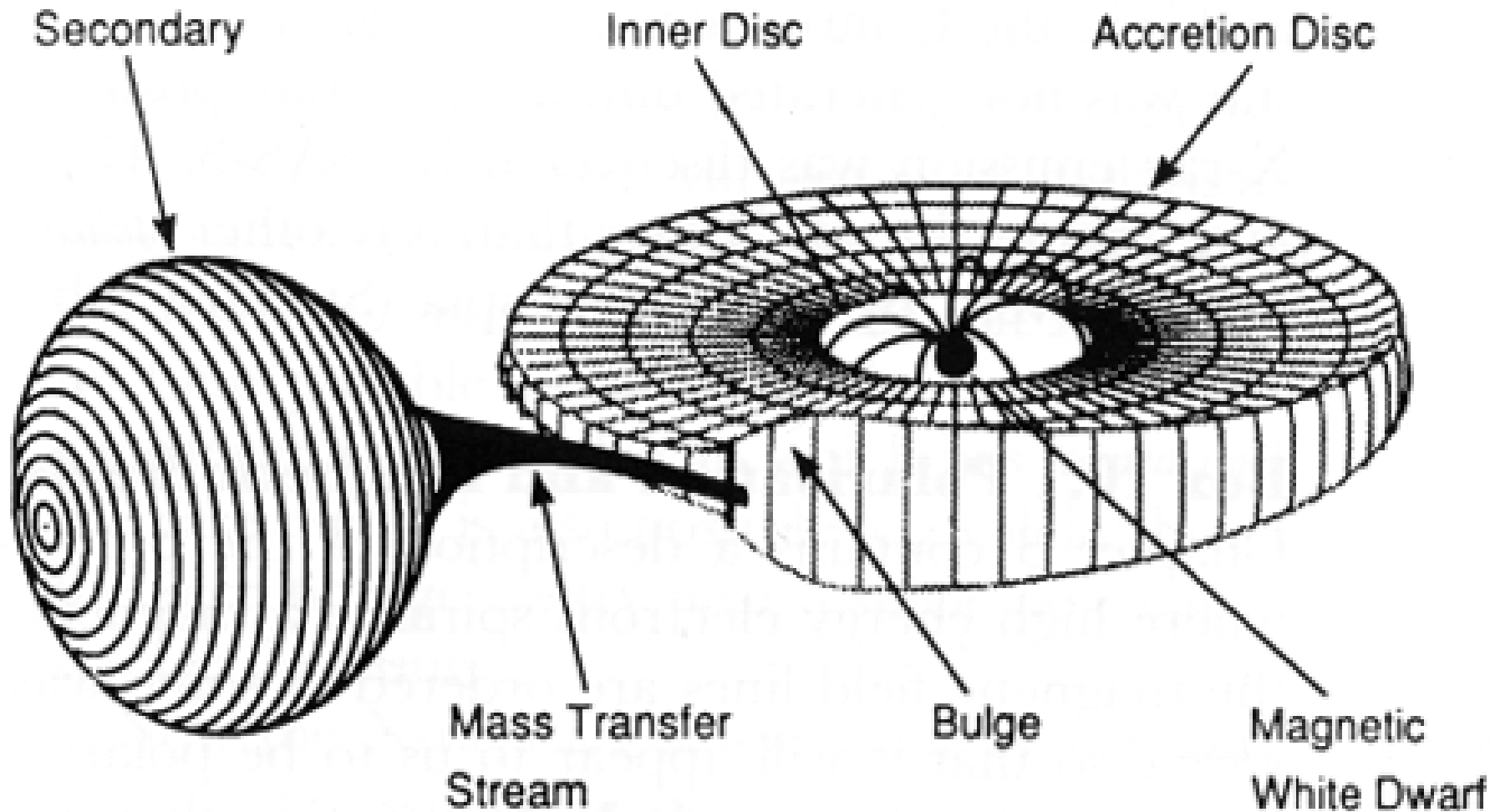
H_2 Close-up Image of Elias 29 Late –Stage
Protostar: Dynamical Age ~ 100 yr
(See poster III.2 by *Mary Barsony*).



Large scale
magnetic field
cannot precess
so fast



Disk truncated by stellar magnetic fields:



Cataclysmic variable systems:

Intermediate polar (QD Her): NO JETS

BASIC INGREDIENTS

- 1) The kinetic energy in the inner part of the accretion disk is transferred stochastically, mainly by shock waves, to internal energy on a dynamical time scale.
 - spatiotemporally localized accretion shocks (SPLASH)
- 2) Radiative cooling time is long.
- 3) Energy is transferred to smaller and smaller mass.
- 4) A small fraction of the mass residing in the inner disk takes most of the energy released in stage (1), and escapes as an outflow.

YSOs

(Torbett, M. V. 1984; Torbett, M. V., & Gilden, D. L. 1992)

Following Soker & Regev (1993) and Soker & Lasota (2004):

Two conditions:

(1) Weakly shocked blobs in the transition layer will expand, and disturb the transition layer in such

a way that a strong shock will develop: $\tau_{\gamma\text{-diffusion}} > \tau_{\text{adiabatic}}$

(2) The strongly shocked gas will cool slowly, such that the thermal pressure has time to accelerate the jet's

material: $\tau_{\gamma\text{-diffusion}} \gg \tau_{\text{Ejection}} > \tau_{\text{Dissipation}}$

These lead to a requirement for the minimum value for the mass accretion rate.

$$\dot{M}_{\text{acc}} > 7 \times 10^{-7} \left(\frac{R}{R_o} \right)^{1.2} \left(\frac{\varepsilon}{0.1} \right)^{1.4} \left(\frac{\alpha}{0.1} \right) M_o \text{ yr}^{-1}$$

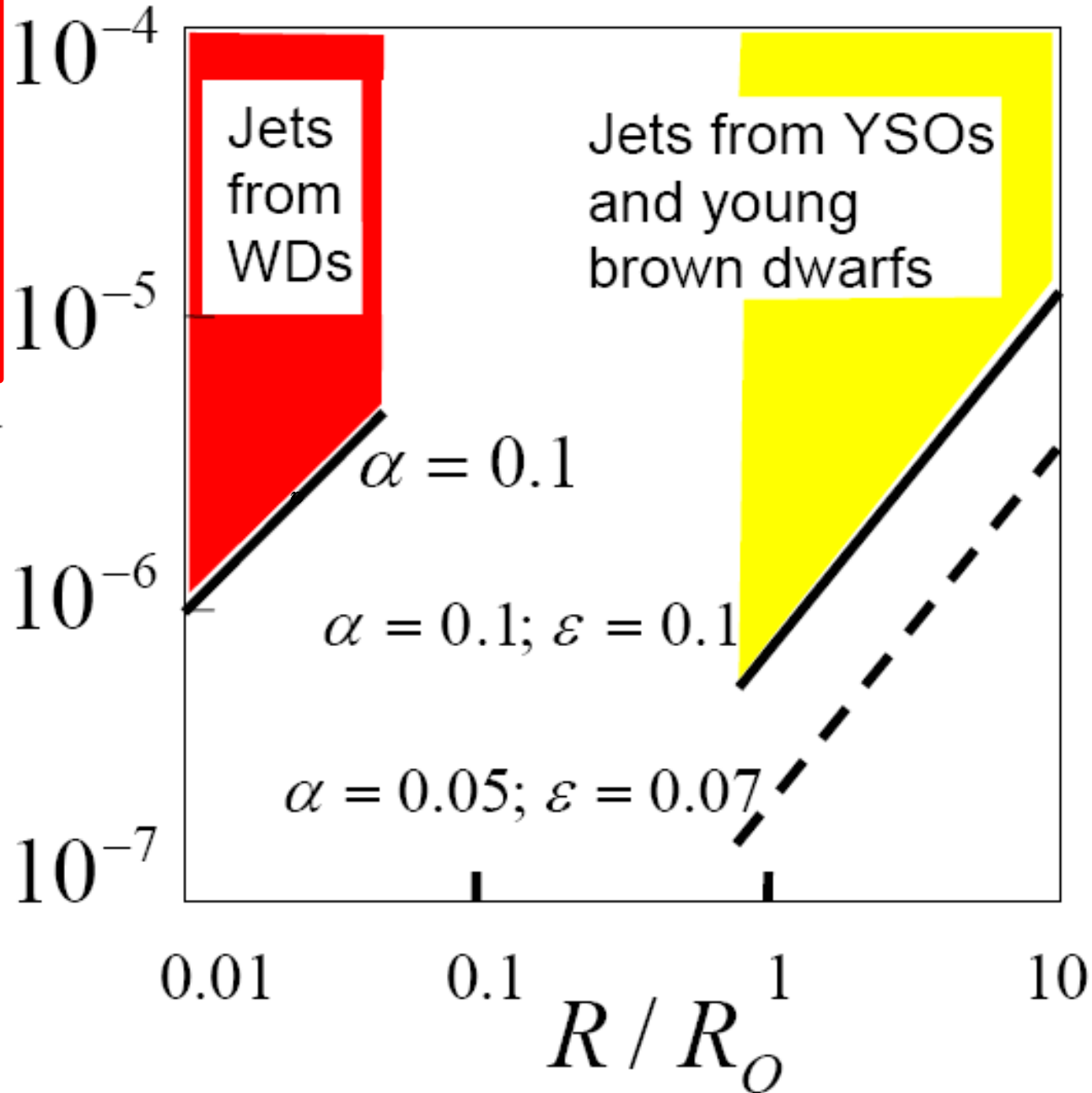
$\varepsilon = \frac{H}{r}$, and α is the usual viscosity parameter.

Accretion into WDs and YSOs

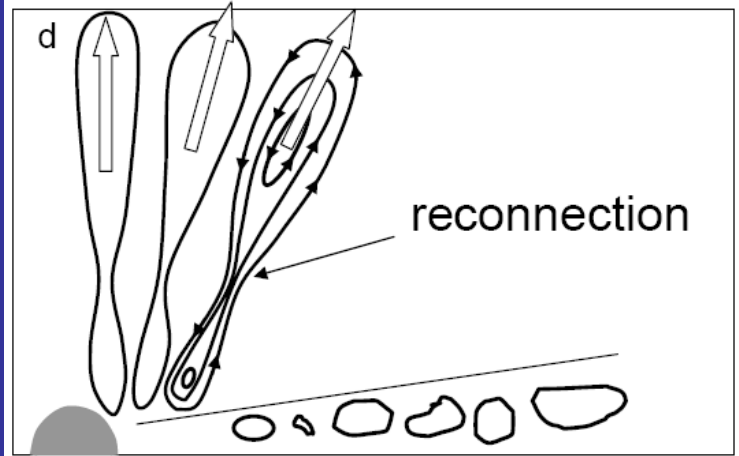
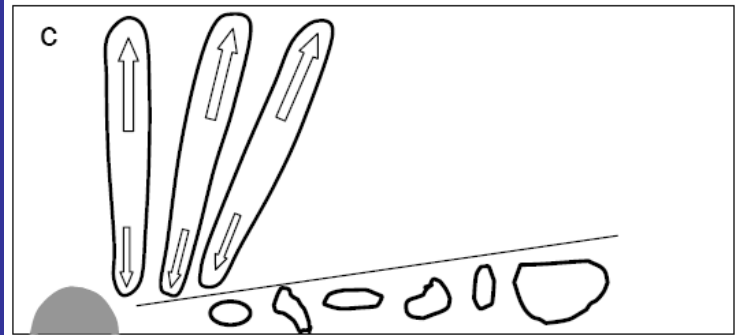
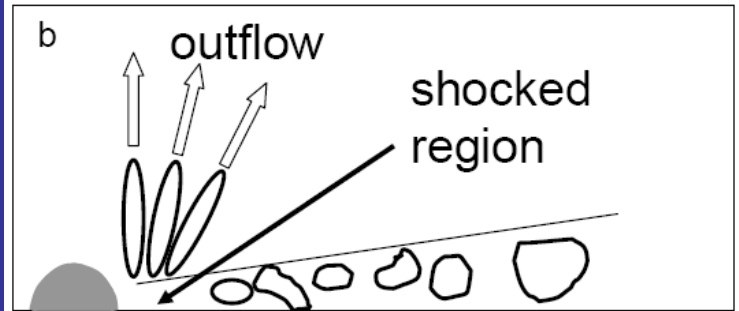
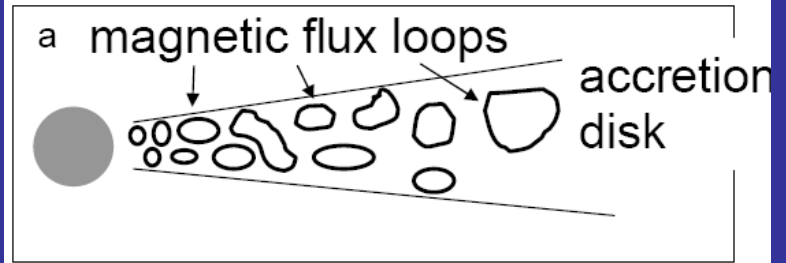
The differences between WDs and YSOs is due to the different opacity dependence on T and density.

$$\dot{M} \propto \frac{1}{K}$$

$$\dot{M}$$



A DOUBLE-STAGE ACCELERATION SCENARIO



Single Stage Acceleration

Kinetic energy of rotating disk's material

Stochastic Shock waves

Internal (e.g., thermal) energy

Pressure gradient

High velocity outflow (jets)

Double Stage Acceleration

Kinetic energy of rotating disk's material

Stochastic Shock waves

Internal (e.g., thermal) energy

Pressure gradient

Outflow

Stretching magnetic flux loops

Magnetic energy in loops

Magnetic field reconnection

High velocity outflow (jets)

Solar Flares

Kinetic energy of solar rotation and convection

$\alpha\omega$ Dynamo

Magnetic fields in the outer envelope

Buoyancy of magnetic tubes

Outflow

Stretching magnetic flux loops

Magnetic energy in loops

Magnetic field reconnection

High velocity ejection

Accretion into Black Holes

